

Design of Ultra-Narrowband Miniaturized High Temperature Superconducting Bandpass Filter

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Abstract—This paper proposes a novel clip-shaped meander-line resonator (CSMLR) to realize miniaturized ultra-narrowband (UNB) bandpass filter design. The main advantage is that it can achieve very weak coupling between adjacent resonators with keeping them very close and introduce transmission zeros (TZs). To further demonstrate the feasibility of using this configuration, a six-pole UNB filter with a fractional bandwidth (FWB) of 0.20% at the center frequency of 1915 MHz was designed on double-sided YBCO high temperature superconducting (HTS) thin films with a thickness of 0.5 mm and dielectric constant of 9.8. The measured responses agree rather well with the simulated ones. The measured results show a maximum insertion loss of 0.31 dB and return loss of 15.5 dB in the passband. Two TZs are generated to improve the passband selectivity, which causes the band-edge steepness better than 50 dB/MHz in both transition bands.

1. INTRODUCTION

In the past few years, plenty of efforts have been made for the design of narrowband bandpass filters, which are increasingly demanded for wireless personal communication services (PCS), global system for mobile communications (GSM) and satellite receiver links [1–7]. The ultra-narrowband (UNB) filters based on high temperature superconducting (HTS) films with very low microwave surface resistance and extremely high unloaded Q-factors can possess desirable low insertion loss, sharp selectivity, and high out-of-band rejection, which are greatly useful in the prevention of interference among wireless systems, especially in the systems with similar frequency bands. At the same time, it can also reduce the bandwidth of transition bands to improve the spectrum utilization.

As described in [8], UNB bandpass filters with a fractional bandwidth (FBW) less than 0.20% are still a challenge, because some difficult conditions must be satisfied to realize a UNB filter. Firstly, a very high unloaded Q-factors value is needed to reduce the loss to an acceptable level as it is well known that the narrower the bandwidth is, the higher the loss would be. Secondly, the UNB filter with extremely weak coupling between adjacent resonators is difficult to achieve within a limited layout area practically, because in order to decrease the coupling strength, increasing the distance between adjacent resonators is a frequently-used method, which is not beneficial for miniaturization. Thirdly, the extremely weak coupling between adjacent resonators would also magnify the influence of unwanted parasitic coupling. So there are a few published papers reporting filters with FBW less than 0.2%. In [9], a five-pole HTS microstrip filter shows a possible FBW of 0.014% at the center frequency of 700 MHz, but there is no transmission zero (TZ) located in transition bands. In [8], an HTS UNB bandpass filter with an FBW of 0.02% in C band is designed, but the measured minimum insertion loss is 3.8 dB. A two-pole HTS filter with FBW of 0.02% for the wireless industry applications is presented in [10], but the return loss and out-of-band rejection are not satisfying.

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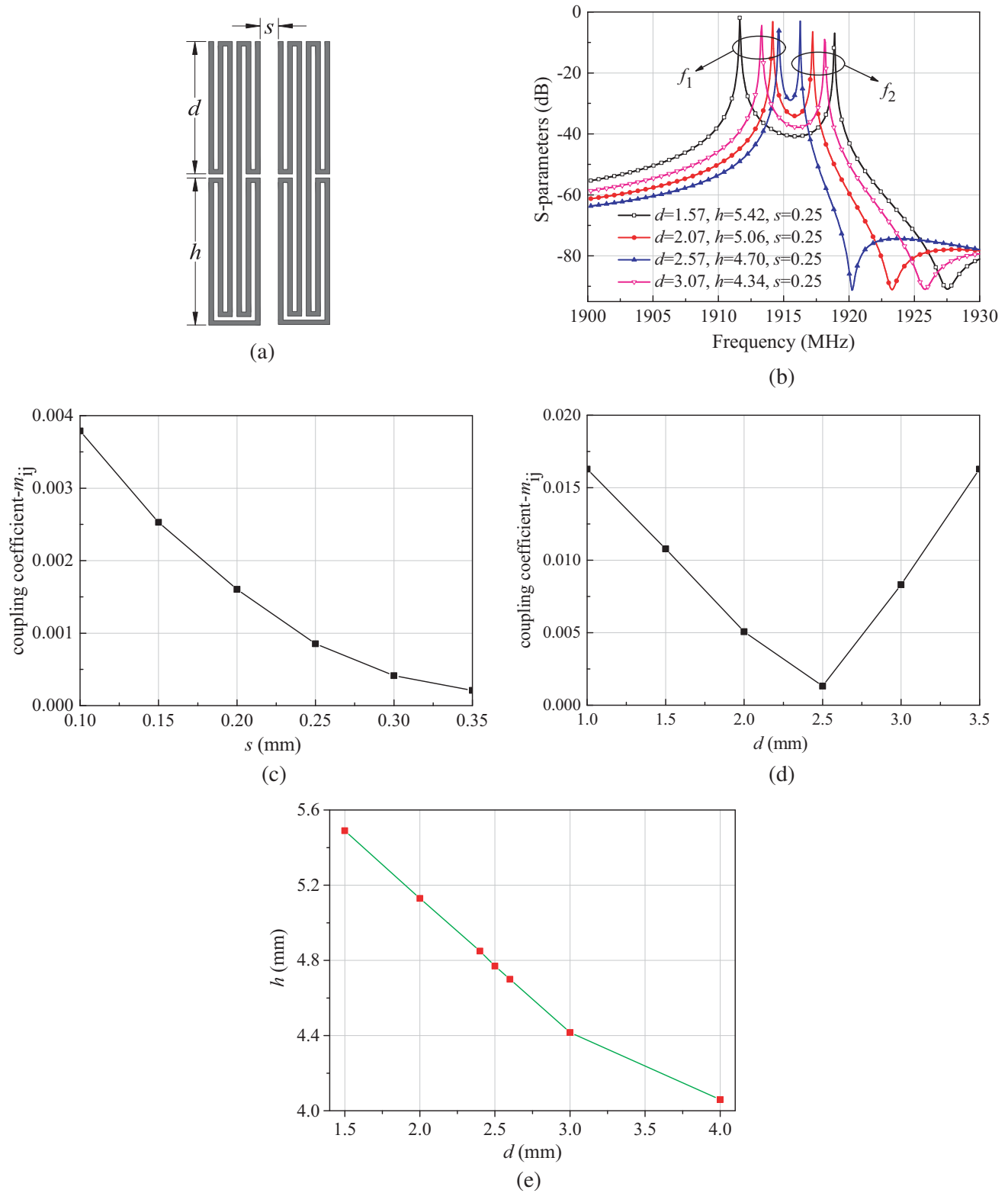


Figure 1. Proposed CSMLR configuration. (a) Layout of two CSMLRs. (b) Relationship between coupling strength and d with $s = 0.25$ mm. (c) Relationship between coupling coefficient and s . (d) Relationship between coupling coefficient and d . (e) Relationship between d and h with f_0 unchanged.

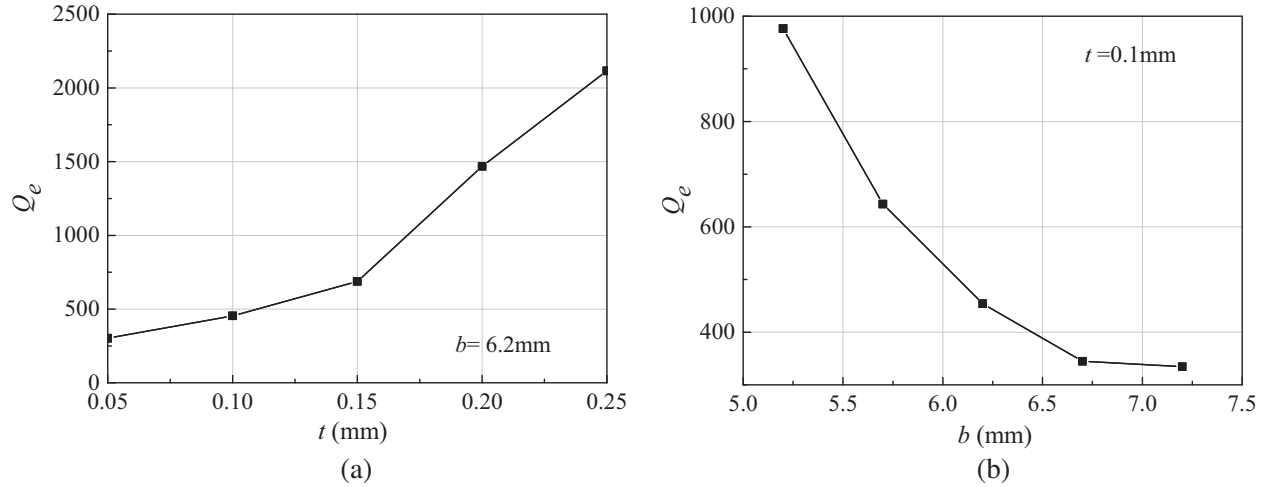


Figure 2. Relationship between the external quality factor Q_e and t or b . (a) t with $b = 6.2$ mm. (b) b with $t = 0.1$ mm.

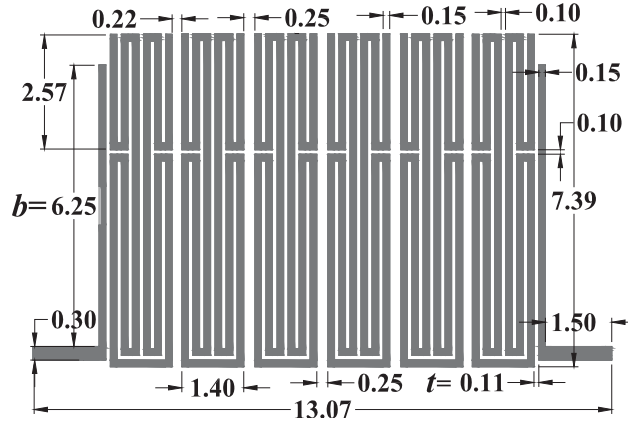


Figure 3. Layout of the proposed six-pole UNB filter (unit: millimeter).

by tuning s in a small range, required coupling coefficients can be obtained without influencing other ones, and compact filter size would also be guaranteed.

Indirect coupling ports are adopted to flexibly adjust the value of external quality factor for achieving the suitable external coupling strength. Figure 2(a) and Figure 2(b) show the simulated external quality factor Q_e as a function of the distance t between the input (output) port and the first (last) resonator, and a function of the height b of ports, respectively. It can be seen that as t increases, Q_e increases to a huge value very quickly, and with the increase of b , Q_e drops to a small value rapidly. So by repeatedly optimizing the distance t and height b , correct external coupling strength can be achieved. Finally, when $t = 0.11$ mm, $b = 6.25$ mm, Q_e is equal to 456.7.

Through the optimization of the full-wave electromagnetic simulation software IE3D, the final associated dimensions of the filter can be determined, and they are marked in Figure 3. It can be observed from Figure 4(b) that a pair of TZs is introduced, however no cross-coupling line is loaded between non-adjacent resonators, or extracted-pole technology is adopted. According to Reference [14], due to the controllable electric and magnetic mixed coupling between adjacent CSMLRs, the TZs (TZ₁ and TZ₂) are introduced. In [14], the authors designed some equivalent circuits of bandpass filters to investigate the electric and magnetic mixed coupling and gave Equations (3), (4) and (5) to present the relationship among inductive/capacitive coupling coefficient (M_C/E_C), odd- and even-mode resonant

frequencies (ω_{od} and ω_{ev}) and the TZ frequency (ω_m).

$$M_C = \frac{\omega_{od}^2 - \omega_{ev}^2}{\omega_{od}^2 + \omega_{ev}^2 - 2\omega_m^2} \quad (3)$$

$$E_C = \frac{\omega_m^2 (\omega_{od}^2 - \omega_{ev}^2)}{2\omega_{od}^2 \omega_{ev}^2 - \omega_m^2 (\omega_{od}^2 + \omega_{ev}^2)} \quad (4)$$

$$m_{ij} = M_C - E_C = \frac{\omega_{od}^2 - \omega_{ev}^2}{\omega_{od}^2 + \omega_{ev}^2} \quad (5)$$

On the basis of the definition of electric and magnetic mixed coupling, Equation (6) is presented below [14]. It can be found that the ratio of the inductive and capacitive coupling coefficients determines the relative position of the TZs. The closer the capacitive coupling and inductive coupling are, the closer the TZs are to the self-resonant frequency (ω_0). Equation (7) can be used to more obviously explain the relationship between TZs and the ratio of the inductive and capacitive coupling coefficients through the simplification of Equation (6) [14].

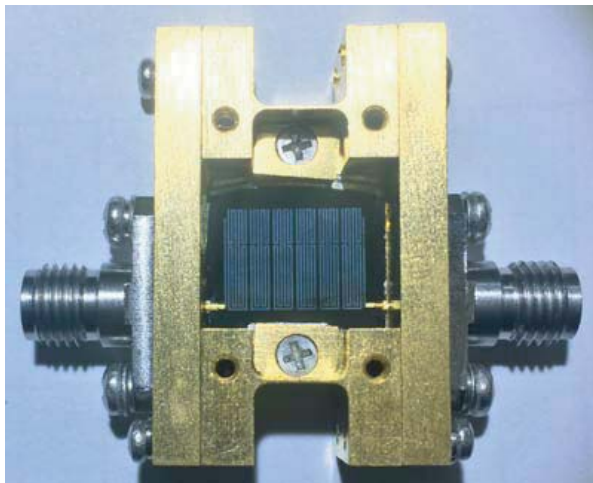
$$\frac{M_C}{E_C} = \frac{L_m/L}{C/C_m} = \frac{L_m C_m}{LC} = \frac{\omega_0^2}{\omega_m^2} \quad (6)$$

$$f_m = f_0 \sqrt{\frac{E_C}{M_C}} \quad (7)$$

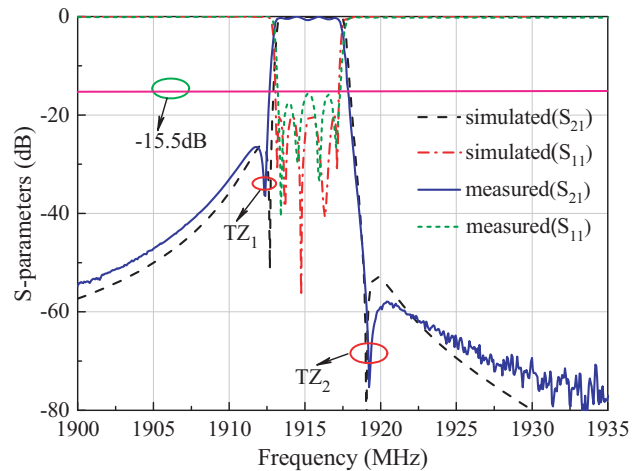
According to Equation (5), when E_C is close to M_C in proposed coupled CSMLRs, and the total coupling m_{ij} would be minimum, which is suitable for UNB design. According to Equation (7), when M_C/E_C in coupled CSMLRs is less than 1, TZ will above ω_0 . Because M_C is very close to E_C in the coupled CSMLRs ($m_{ij} \approx 0$), M_C/E_C less or more than 1 can be easily realized by adjusting the parameter d slightly in Figure 1(a), so that the control of TZs can be achieved in the proposed UNB filter design.

3. FABRICATION AND MEASUREMENT OF UNB FILTER

The UNB filter was fabricated on a 0.5-mm thick MgO wafer with double-sided YBCO HTS thin films. The relative dielectric constant of MgO is 9.8. One side of the HTS films is patterned into the filter



(a)



(b)

Figure 4. Photograph and results of fabricated UNB filter. (a) Photograph of fabricated six-pole HTS UNB filter. (b) Measured and simulated S -parameters of the filter.

circuit by the standard procedure of photolithography, the other side is used for grounding. The eventual circuit size is $13.07 \text{ mm} \times 7.39 \text{ mm}$ ($0.26\lambda_g \times 0.15\lambda_g$, λ_g is the wavelength in the dielectric).

A physical photograph of the fabricated HTS UNB filter is displayed in Figure 4(a). The packaged HTS UNB filter was cooled down to a temperature of 77 K in a cryogenic cooler and measured by HP8753 network analyzer, after the full two-port calibration for reflection and transmission measurements was performed at ambient temperature, and Figure 4(b) illustrates the simulated and measured frequency responses. It can be observed that the FWB is 0.20% at the center frequency of 1915 MHz, the same as the simulated one. And due to the existence of TZs, the passband selectivity is improved effectively. The band-edge steepness reaches over 50 dB/MHz in both transition bands, and the out-of-band rejection is better than 70 dB at the frequency of 1919 MHz. The maximum in-band insertion loss is 0.31 dB with corresponding unloaded Q-factors about 67000. The return loss is better than 15.5 dB within the passband. In general, the measured results are in good agreement with simulated ones. Table 2 is the comparison between the proposed UNB filter with the reported ones.

Table 2. Performance comparison with reported UNB filters.

Ref.	f (MHz)	FBW (%)	Q-factors	VSWR	TZs	Pole	Size ($\lambda_g \times \lambda_g$)
[6]	9900	0.65	24000	1.58	No	6	1.81×0.27
[11]	900	0.27	10800	1.20	No	5	No
[12]	1968	0.25	48000	1.38	Yes	8	0.78×0.27
[13]	1967.5	0.76	17000	1.9	Yes	18	1.23×0.27
This work	1915	0.2	67000	1.40	Yes	6	0.26×0.15

4. CONCLUSION

A novel CSMLR is proposed to reach very weak coupling between adjacent resonators for miniaturized UNB filter design. A pair of TZs can be produced to realize sharp passband selectivity in the design of the six-pole HTS UNB filter. The simulated six-pole UNB filter was fabricated finally to verify the correctness and feasibility of this structure. The experimental results show the filter possesses high performance, which can be applied to modern wireless communication systems.

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