

## MULTICONDUCTOR REDUCTION METHOD FOR MODELING CROSSTALK OF COMPLEX CABLE BUNDLES IN THE VICINITY OF A 60 DEGREE CORNER

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**Abstract**—This paper presents a multiconductor reduction method for modeling electromagnetic crosstalk of complex cable bundles in the vicinity of a 60 degree corner. Based on the image theory and wide separation assumption, the per-unit-length parameters of the cable bundle can be obtained analytically. A modified six-step procedure is established to define the electrical and geometrical characteristics of the reduced cable bundle model compared with the original equivalent cable bundle method (ECBM). Numerical simulations are performed to demonstrate the viability and effectiveness of the method. This work can find wide applications in real environments.

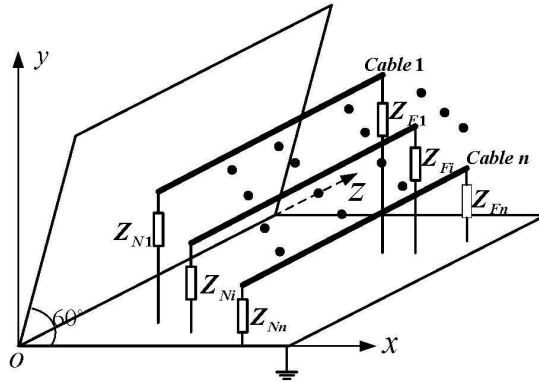
### 1. INTRODUCTION

With the development of electronic technology, electromagnetic (EM) environment [1–3] of electronic devices in aeronautic and automotive systems becomes more and more complex. Much electromagnetic interference (EMI) and electromagnetic compatibility (EMC) problems associated with cable bundles [4, 5] connecting these devices should be taken into consideration. In recent years, a new simplification method called ECBM [6] has been proposed which is based on the main assumption that the common-mode response is more critical than the differential-mode response when considering external EM waves coupling to cables. This effective method has been applied in the EM

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*Received 19 December 2012, Accepted 17 January 2013, Scheduled 23 January 2013*

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**Figure 1.** Illustration of  $n$ -conductor transmission lines in the vicinity of a 60 degree corner.

emissions [7, 8] and crosstalk [9, 10] of complex cable bundles over a large frequency range with some specific adjustments. It has also been adapted to predict the crosstalk of a cable bundle in the cylindrical cavity which is considered as the ground return [11].

However, in real circumstances, cable bundles are commonly set in more complicated environments. Cable bundles in the vicinity of a corner is one of the most common environments. In practice, there are many articles about the application of transmission line [12–15]. And in this paper, we put forward a multiconductor reduction method for modeling EM crosstalk [16] of a complex cable bundle in the vicinity of a 60 degree corner. Fig. 1 shows the geometry of an  $n$ -conductor transmission lines (TL) in the vicinity of a 60 degree corner laid along the  $z$  axis and parallel to the  $xoz$  and  $yo z$  plane which are both considered as PEC ground planes, and the angle between them is 60 degrees.  $Z_{N_i(F_i)}$  ( $i = 1, 2, \dots, n$ ) represent the near (far) end termination loads of the  $n$ -conductor TL. According to the multiconductor transmission line (MTL) and image theory [17], we can obtain the formula of the per-unit-length (p.u.l.) parameters of the cable bundle to make the ECBM fit for this new environment. Different from [11], a modified six-step procedure is established to simplify the EM crosstalk problem in this situation.

The organization of this paper is as follows. In Section 2, a modified six-step procedure is given, and the derivation and validation of the p.u.l. parameters of the cable bundle is given in detail. In Section 3, simulation results are given to validate the method, and in the final section, some comments on the proposed method are presented.

## 2. DERIVATION AND VALIDATION OF THE P.U.L. PARAMETERS OF MTL IN THE VICINITY OF A 60 DEGREE CORNER

### 2.1. Derivation of the p.u.l. Parameters

For the derivation, the following assumptions are necessary [17, 18]:

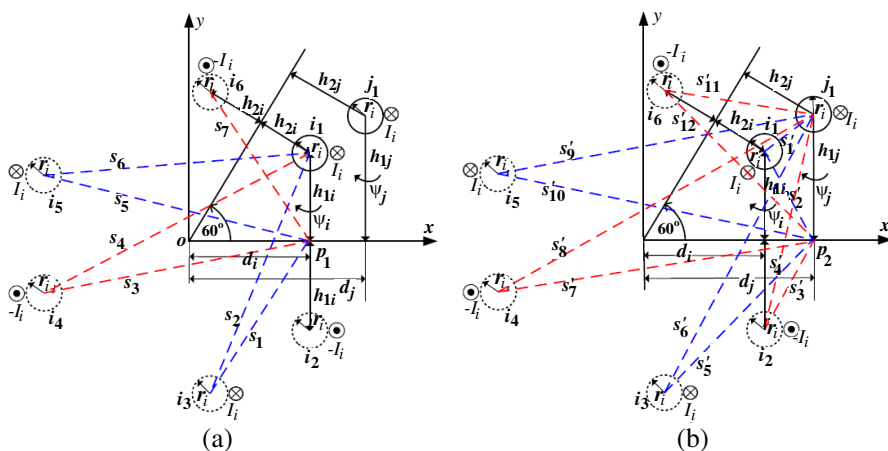
*i*) all conductors are PEC and the medium surrounding them is lossless;

*ii*) only transverse electromagnetic (TEM) mode is propagated along the conductor;

*iii*) the interval between each two conductors is wide enough that the charge and current distribution around each conductor can be considered as uniform.

According to the image theory [17], we can get the formula of the p.u.l. self-inductance  $L_{ii}$  of the  $i$ th conductor and mutual-inductance  $L_{ij}$  between the  $i$ th and  $j$ th conductors. Fig. 2 shows the derivation of  $L_{ii}$  and  $L_{ij}$  on the basis of the image theory.

$$\begin{aligned}
 L_{ii} &= \frac{\psi_i}{I_i} \Big|_{I_k=0} (k = 1 \dots n, k \neq i) = \frac{\mu_0}{2\pi} \ln \left( \frac{h_{1i}}{r_i} \right) + \frac{\mu_0}{2\pi} \ln \left( \frac{2h_{1i}}{h_{1i}} \right) \\
 &\quad - \frac{\mu_0}{2\pi} \ln \left( \frac{s_2}{s_1} \right) + \frac{\mu_0}{2\pi} \ln \left( \frac{s_4}{s_3} \right) + \frac{\mu_0}{2\pi} \ln \left( \frac{s_6}{s_5} \right) - \frac{\mu_0}{2\pi} \ln \left( \frac{s_7}{2h_{2i}} \right) \\
 &= \frac{\mu_0}{2\pi} \ln \left( \frac{4h_{1i} \cdot h_{2i} \cdot s_1 \cdot s_4 \cdot s_6}{r_i \cdot s_2 \cdot s_3 \cdot s_5 \cdot s_7} \right), \tag{1}
 \end{aligned}$$



**Figure 2.** (a) The p.u.l. self-inductance  $L_{ii}$  of the  $i$ th conductor. (b) The p.u.l. mutual-inductance  $L_{ij}$  between the  $i$ th and  $j$ th conductors.

in which  $s_1 = s_7$ ,  $s_3 = s_6$ ,  $h_{2i} = \frac{\sqrt{3}d_i - h_{1i}}{2}$  and

$$s_2 = \sqrt{(2h_{1i} + h_{2i})^2 + 3h_{2i}^2}, \quad (2)$$

$$s_4 = \sqrt{\left(\sqrt{3}h_{1i} + \sqrt{3}h_{2i}\right)^2 + (h_{1i} + h_{2i})^2}, \quad (3)$$

$$s_5 = \sqrt{\left(\sqrt{3}h_{1i} + \sqrt{3}h_{2i}\right)^2 + (h_{1i} - h_{2i})^2}, \quad (4)$$

in which  $r_{i(j)}$  is the radius of the  $i(j)$ th conductor. conductors  $i_2$ ,  $i_3$ ,  $i_4$ ,  $i_5$ ,  $i_6$  are the image conductors of the  $i$ th conductor.  $I_{i(j)}$ ,  $-I_{i(j)}$  are the currents flow along the  $i(j)$ th conductor and its image conductors.  $\psi_i$  is the magnetic flux which passes through the reference plane.  $h_{1i(j)}$  represents the vertical distance between the conductor  $i_1(j_1)$  and the  $xoz$  plane.  $h_{2i(j)}$  represents the vertical distance between the conductor  $i_1(j_1)$  and the inclined plane.  $d_i$ ,  $d_j$  are the distances between the conductors  $i_1$ ,  $j_1$  and  $yo$ z plane, respectively.  $s_1$ ,  $s_3$ ,  $s_5$ ,  $s_7$  are the distances between the image conductors  $i_2$ ,  $i_3$ ,  $i_4$ ,  $i_5$ ,  $i_6$  and the reference point  $p_1$ , respectively.  $s_2$ ,  $s_4$ ,  $s_6$  are the distances between the image conductors  $i_3$ ,  $i_4$ ,  $i_5$  and conductor  $i_1$ , respectively. The formula of the mutual-inductance can be obtained as follows,

$$L_{ij} = \frac{\psi_i}{I_j} |_{I_k=0} (k=1 \dots n, k \neq j) = \frac{\mu_0}{2\pi} \ln\left(\frac{s'_2}{s'_1}\right) + \frac{\mu_0}{2\pi} \ln\left(\frac{s'_4}{s'_3}\right) - \frac{\mu_0}{2\pi} \ln\left(\frac{s'_6}{s'_5}\right) + \frac{\mu_0}{2\pi} \ln\left(\frac{s'_8}{s'_7}\right) + \frac{\mu_0}{2\pi} \ln\left(\frac{s'_{10}}{s'_9}\right) - \frac{\mu_0}{2\pi} \ln\left(\frac{s'_{12}}{s'_{11}}\right) = \frac{\mu_0}{2\pi} \ln\left(\frac{s'_2 \cdot s'_4 \cdot s'_5 \cdot s'_8 \cdot s'_{10} \cdot s'_{11}}{s'_1 \cdot s'_3 \cdot s'_6 \cdot s'_7 \cdot s'_9 \cdot s'_{12}}\right), \quad (5)$$

in which  $s'_2 = s'_3$ ,  $s'_5 = s'_{12}$ ,  $s'_7 = s'_{10}$ ,  $d_1 = d_j - d_i$  and

$$h_{2j} = \frac{\sqrt{3}d_j - h_{1j}}{2}, \quad (6)$$

$$s'_1 = \sqrt{(h_{1j} - h_{1i})^2 + d_1^2}, \quad (7)$$

$$s'_4 = \sqrt{(h_{1j} + h_{1i})^2 + d_1^2}, \quad (8)$$

$$s'_6 = \sqrt{(h_{1i} + h_{2i} + h_{1j})^2 + (d_1 + \sqrt{3}h_{2i})^2}, \quad (9)$$

$$s'_8 = \sqrt{(h_{1j} + h_{2i})^2 + (d_1 + \sqrt{3}h_{1i} + \sqrt{3}h_{2i})^2}, \quad (10)$$

$$s'_9 = \sqrt{(h_{1j} - h_{2i})^2 + (d_1 + \sqrt{3}h_{1i} + \sqrt{3}h_{2i})^2}, \quad (11)$$

$$s'_{11} = \sqrt{(h_{1j} - h_{2i} - h_{1i})^2 + (d_1 + \sqrt{3}h_{2i})^2}, \quad (12)$$

in which  $s'_2, s'_3, s'_5, s'_7, s'_{10}, s'_{12}$  are the distances between the conductors  $i_1, i_2, i_3, i_4, i_5, i_6$  and the reference point  $p_2$ , respectively.  $s'_1, s'_4, s'_6, s'_8, s'_9, s'_{11}$  are the distances between the image conductors  $i_1, i_2, i_3, i_4, i_5, i_6$  and conductor  $j_1$ , respectively.

## 2.2. Validation of the p.u.l. Parameters

In order to verify the validity of (2) and (6), we compare the results calculated from the analytical expressions with the results obtained by numerical simulations through some examples.  $[L_N]$  represents the numerical results and  $[L_A]$  represents the analytical one.

1)  $r_i = r_j = 0.5$  mm,  $h_{1i} = 40$  mm,  $h_{1j} = 43$  mm,  $d_i = 50$  mm,  $d_j = 51$  mm and the p.u.l. inductance matrices (in nanohenry/meter)

$$[L_N] = \begin{bmatrix} 872.3 & 501.0 \\ & 872.5 \end{bmatrix}, \quad [L_A] = \begin{bmatrix} 868.1 & 505.1 \\ & 868.0 \end{bmatrix}. \quad (13)$$

2)  $r_i = 0.5$  mm,  $r_j = 0.6$  mm,  $h_{1i} = 41$  mm,  $h_{1j} = 43$  mm,  $d_i = 50$  mm,  $d_j = 52$  mm and the p.u.l. inductance matrices (in nanohenry/meter)

$$[L_N] = \begin{bmatrix} 867.1 & 524.9 \\ & 838.9 \end{bmatrix}, \quad [L_A] = \begin{bmatrix} 866.4 & 523.8 \\ & 837.1 \end{bmatrix}. \quad (14)$$

Other comparison results which show excellent agreements and not reported here fully demonstrate that the analytical expressions which is based on the wide separation assumption can be accepted and introduced in the ECBM to define the electrical and geometrical characteristics of the reduced cable bundle model in the vicinity of a 60 degree corner.

## 3. PRESENTATION OF THE REDUCTION METHOD

In this section, a modified six-step procedure is established to simplify the EM crosstalk problem of complex cable bundles in the vicinity of a 60 degree corner. Compared with [11], Steps I, II, IV, V are identical. Thus, these steps are omitted only to avoid repetition. Only Step III is quite different and reported in detail.

### *Step III: Reduced Cable Bundle Cross-Section Geometry*

This step is to build the reduced cable model cross-section geometry. It is realized thanks to the knowledge of the  $[L_{reduced}]$  and  $[C_{reduced}]$  matrices. As far as this new environment is concerned, a new optimization process made of six phase is necessary.

1) Phase 1: Estimate the height  $h_{1i}$  and  $h_{2i}$  above the ground plane of each equivalent conductor.  $h_{1i}$  and  $h_{2i}$  correspond to the average heights of all the conductors in group  $i$  to  $xoz$  and the inclined plane.

2) Phase 2: Calculate the radius  $r_i$  of each equivalent cable according to (1)

$$r_i = \frac{2 \cdot h_{1i} \cdot h_{2i} \cdot s_4}{\exp\left(\frac{2\pi L_{ii\_reduced}}{\mu_0}\right) \cdot s_2 \cdot s_5}. \quad (15)$$

3) Phase 3: Calculate the distance  $d_{ij}$  between each two equivalent cables according to (5)

$$d_{ij} = \frac{s'_4 \cdot s'_8 \cdot s'_{11}}{\exp\left(\frac{2\pi L_{ij\_reduced}}{\mu_0}\right) \cdot s'_6 \cdot s'_9}. \quad (16)$$

4) Phase 4: Adjust  $r_i$ ,  $d_{ij}$  determined by the above procedures using a dichotomic optimization realized with exact electrostatic calculations in the error range. This phase allows a reduction of the first-estimate errors involved in (15) and (16).

5) Phase 5: Determine the thickness of the dielectric coating surrounding the conductor of each equivalent cable while avoiding dielectric coating overlapping [6].

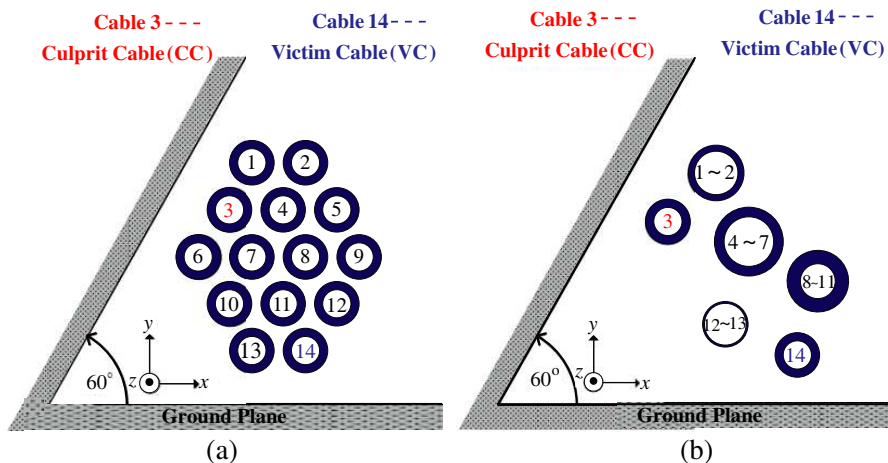
6) Phase 6: Calculate the relative permittivity  $\varepsilon_r$  of each cable dielectric coating according to the  $[C_{reduced}]$  matrix using an electrostatic calculation [6].

#### 4. VALIDATION OF THE REDUCTION METHOD FOR CROSSTALK PREDICTION THROUGH NUMERICAL SIMULATIONS

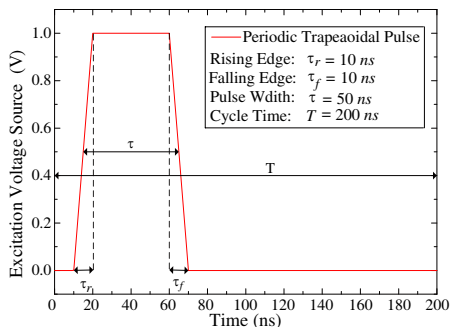
In this section, a 14-conductor point-to-point connected cable bundle, 1 m long, set in the vicinity of a 60 degree corner shown in Fig. 3(a) is investigated, in which all cables are single wire cables with the radius of 0.5 mm and surrounded by a dielectric coating with the thickness

**Table 1.** Termination loads of the 14-complete cable bundle model (unit:  $\Omega$ ).

Conductor	1	2	3	4	5	6	7
Near End	50	60	50	1.5k	1.2k	900	1.2k
Far End	40	28	50	1k	1k	2k	1.5k
Conductor	8	9	10	11	12	13	14
Near End	2.3k	1.8k	900	1.2k	55	38	50
Far End	25	30	70	48	1.6k	1.5k	100



**Figure 3.** Cross-section geometry of (a) the complete model and (b) the reduced model.



**Figure 4.** The trapezoidal pulse waveform of the voltage source excited on Cable 3.

of 0.3 mm and dielectric constant of  $\epsilon_r = 2.5$  and  $\mu_r = 1.0$ . The distance between the two neighboring lines is 3 mm horizontally and vertically. The near end of Cable 3 (culprit cable) is excited with a periodic trapezoidal pulse voltage source shown in Fig. 4. Cable 14 serves as the victim cable. The p.u.l. parameter matrices inductance  $[L]$  (in nanohenry/meter) and capacitance  $[C]$  (in picofrad/meter) of the cable bundle are listed in (17) and (18). Meanwhile, in order to make the problem simpler, we only consider real loads at the two terminals of the cable bundle, which are listed in Table 1. The common-mode characteristic impedance  $Z_{mc}$  which can be determined by modal analysis [8] equals  $90 \Omega$ . According to the grouping process in Section 2, the conductors of the cable bundle can be sorted into four

groups as follows. 1) group 1: Cables 1 ~ 2; 2) group 2: Cables 4 ~ 7; 3) group 3: Cables 8 ~ 11; 4) group 4: Cables 12 ~ 13.

The p.u.l. parameter matrices  $[L]$  (in nanohenry/meter) and capacitance  $[C]$  (in picofrad/meter) of the reduced cable bundle model can be easily obtained as follows [8],

$$[L] = \begin{bmatrix} 725 & 339 & 318 & 256 & 175 & 218 & 219 & 197 & 153 & 170 & 170 & 131 \\ & 786 & 348 & 371 & 164 & 226 & 266 & 264 & 163 & 196 & 212 & 150 \\ & & 783 & 392 & 239 & 349 & 371 & 304 & 230 & 269 & 267 & 200 \\ & & & 833 & 193 & 287 & 395 & 415 & 211 & 271 & 309 & 204 \\ & & & & 735 & 346 & 244 & 192 & 328 & 263 & 211 & 228 \\ & & & & & 786 & 394 & 288 & 355 & 375 & 307 & 274 \\ & & & & & & 831 & 437 & 291 & 397 & 416 & 275 \\ & & & & & & & 875 & 238 & 330 & 436 & 251 \\ & & & & & & & & 796 & 400 & 292 & 283 \\ & & & & & & & & & 836 & 439 & 403 \\ & & & & & & & & & & 876 & 334 \\ & & & & & & & & & & & 837 \end{bmatrix}_{12 \times 12}, \quad (17)$$

$$[C] = \begin{bmatrix} 26.4 & -7.6 & -4.4 & -0.7 & -0.5 & -0.2 & -0.1 & -0.1 & 0 & 0 & 0 & 0 \\ & 25.6 & -4.4 & -6.4 & -0.1 & -0.1 & -0.3 & -1.0 & 0 & 0 & -0.2 & 0 \\ & & 30.6 & -5.8 & -0.5 & -4.1 & -4.2 & -0.1 & 0 & -0.2 & 0 & 0 \\ & & & 27.4 & 0 & -0.4 & -4.3 & -6.4 & 0 & -0.1 & -0.6 & 0 \\ & & & & 26.4 & -5.9 & -0.2 & 0 & -6.0 & -0.5 & 0 & -0.8 \\ & & & & & 30.6 & -5.7 & -0.2 & -4.3 & -4.2 & -0.4 & -0.3 \\ & & & & & & 30.6 & -6.0 & -0.4 & -4.2 & -4.3 & -0.1 \\ & & & & & & & 25.5 & -0.1 & -0.5 & -6.5 & -0.3 \\ & & & & & & & & 27.6 & -5.8 & -0.2 & -6.3 \\ & & & & & & & & & 30.6 & -5.8 & -4.4 \\ & & & & & & & & & & 27.4 & -0.9 \\ & & & & & & & & & & & 25.4 \end{bmatrix}_{12 \times 12} \quad (18)$$

$$[L_{reduced}] = \begin{bmatrix} 547 & 259 & 203 & 166 \\ & 422 & 307 & 250 \\ & & 470 & 362 \\ & & & 598 \end{bmatrix}_{4 \times 4}, \quad (19)$$

$$[C_{reduced}] = \begin{bmatrix} 36.7 & -16.8 & -1.69 & -0.4 \\ & 81.4 & -36.8 & -2.3 \\ & & 80.3 & -28.0 \\ & & & 51.1 \end{bmatrix}_{4 \times 4}. \quad (20)$$

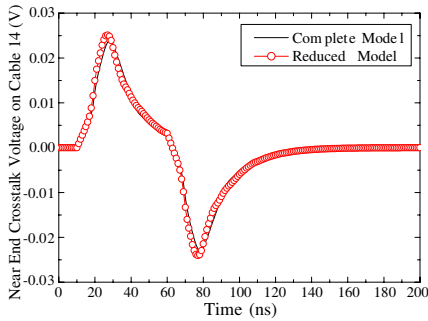
After applying the six procedures described in Section 2, we obtain the cross-section geometry of the reduced cable bundle model composed of six equivalent conductors shown in Fig. 3(b). The equivalent termination loads connected to each end of all conductors and some corresponding parameters of the reduced cable bundle are listed in Table 2.

The crosstalk voltages on the near and far ends of cable 14 can be obtained by applying the MTLN to the complete and reduced cable

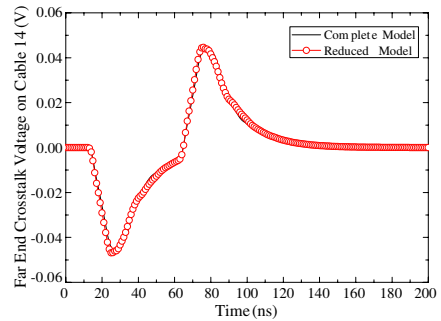
**Table 2.** Termination loads (unit:  $\Omega$ ) and some parameters (unit: mm) of the 6-reduced cable bundle model.

Conductor	1-2	3	4-7	8-11	12-13	14
Near End	27.3	50	290.3	340.7	22.5	50
Far End	16.5	50	292.7	9.2	774.2	100
Conductor Radius	1.5	0.5	2.1	2.2	1.0	0.5
Insulator Thickness	0.3	0.3	0.1	0.2	0.2	0.3





**Figure 5.** Comparison of the near end crosstalk voltage in the time domain on Cable 14 between the complete and reduced cable bundle models.



**Figure 6.** Comparison of the far end crosstalk voltage in the time domain on Cable 14 between the complete and reduced cable bundle models.

bundle models and are shown in Figs. 5 and 6. The comparison gives excellent agreement.

## 5. CONCLUSIONS

This paper details a modified multiconductor reduction method for a cable bundle in the vicinity of a 60 degree corner and presents a modified six-step procedure to simplify the EM crosstalk problem. The excellent agreement validates the efficiency and the advantages of the method.

In this numerical simulation, the total computation time is reduced by a factor of 4.1 (complete model costs 53 seconds, reduced model costs 13 seconds) after equivalence of the complete model by using the method of MTLN theory, which have been performed on a 2.1-GHz processor and a 2.0-GB RAM memory computer. From these results, we can show that this method can significantly reduce the prediction time and memory requirements. It could be expected that with the cable number in the original cable bundle increases, we can cut down much more computation time and memory.

## ACKNOWLEDGMENT

This work was supported by the National Natural Science Foundation of China for Young Scholars under Grant No. 61102033, the Fundamental Research Funds for the Central Universities under Grant NS2012096, the Foundation of State Key Laboratory of Millimeter

Waves, Southeast University, P. R. China under Grant No. K201302 and the Aeronautical Science Foundation of China under grant No. 20128052063.

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