

HIGH-FREQUENCY METHOD ANALYSIS ON SCATTERING FROM HOMOGENOUS DIELECTRIC OBJECTS WITH ELECTRICALLY LARGE SIZE IN HALF SPACE

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Abstract—The high-frequency method for solving the scattering from homogeneous dielectric objects with electrically large size in half space is presented in this paper. In order to consider the scattering fields of the targets in half space, the half-space physical optics method is deduced by introducing the half-space Green's function into the conventional physical optics method (PO). Combined with the graphical-electromagnetic computing method to read the geometry information of all visible facets, the equivalent currents and the reflection coefficients are utilized to account of the homogenous dielectric objects with half-space physical optics method in half space. The numerical results show that this method is efficient and accurate.

1. INTRODUCTION

With the development of the electromagnetic, the researchers have directed significant attention toward the targets in half space, especially the electrically large dielectric objects above the earth or the sea, such as geophysical, remote sensing, wave propagation. A variety of the RCS prediction software have been widely used for interactive modeling, design, and analysis of aircraft with RCS specifications in free space, such as XPATCH, CADDSCAT, RESPECT, GRECO and RANUS [1], which are based on the high-frequency techniques [2–7] and provide an efficient tool to obtain real-time results. But in the scope of author's knowledge, it is not reported that the homogeneous dielectric object with electrically large size in half space was calculated by the high-frequency method.

Numerous authors have derived a variety of methods which are used for computing the high-frequency radar cross section (RCS) of complex radar targets in half space. The scattering fields of the targets in half space could be computed by using the Finite element method (FEM) [8], Finite-difference time-domain methods (FDTD) [9], but they cause some difficulties in the solution procedure [12]. Some researchers try to combine the half-space Green's function and the Method of moment (MOM) to consider the electrically large targets in half space [10–12], but it also requires more mass storage memory, which results in lower processing speed and computational efficiency. In this paper, based on the analysis of the electrically large conductive targets in half space [13], we have further introduced the high-frequency method into the homogeneous dielectric object with electrically large size in half space.

This paper focuses on the prediction of radar cross section of homogeneous dielectric object with electrically large size in half space. In order to consider the scattering fields of the electrically large target in half space, the half-space physical optics method (PO) for dielectric object is obtained by introducing the half-space Green's function into the PO method for dielectric object. Combined with the graphical-electromagnetic computing (GRECO) method, the shadow regions are eliminated by displaying lists technology of OpenGL to rebuild the target, and the geometry information is attained by reading the color and depths of each pixel [4]. The results show that the half-space physical optics method for dielectric objects is efficient and accurate. For illustration purposes, we limit our presentation to the monostatic case for homogeneous dielectric objects with electrically large size in half space.

2. THEORY

Consider an arbitrarily shaped dielectric object illuminated by a plane wave in the half space, surface s is assumed to represent a closed surface of target. The far zone scattering electric and magnetic fields of the homogenous medium in half space can be expressed as:

$$E_s(r') = -j\omega A(r') - \nabla V_e - \frac{1}{\varepsilon_0} \nabla \times F, \quad r' \in S \quad (1)$$

$$H_s(r') = -j\omega F(r') - \nabla V_m + \frac{1}{\mu_0} \nabla \times A, \quad r' \in S \quad (2)$$

where A and V_e represent electric vector and scalar potentials, F and V_m represent magnetic vector and scalar potentials.

As show in Fig. 1, in order to solve the scattering fields of the homogenous dielectric object with electrically large size in half space, the half-space Green's function is introduced into the physical optics method to calculate the contributions of scattering between the target and the ground plane.

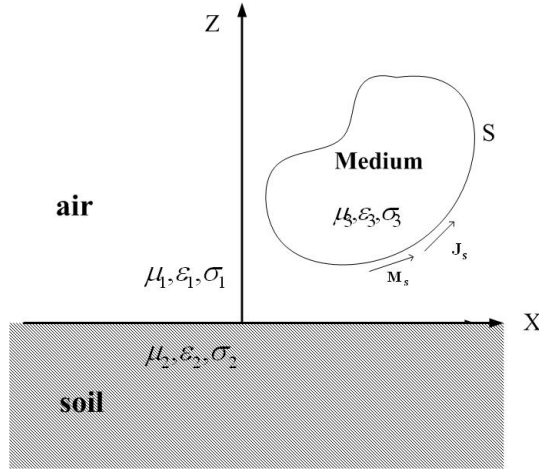


Figure 1. The reflected field of the ground and target in half space.

It is permissible to express the vector and the scalar potentials as follows [14]:

$$A(r) = \mu_0 \int_S \bar{G}_A(r, r') \cdot J(r') dS' \quad (3)$$

$$V_e(r) = \frac{-1}{j\omega\epsilon_0} \int_S \left[G_x^{qe}(r, r') \frac{\partial J_x}{\partial x'} + G_y^{qe}(r, r') \frac{\partial J_y}{\partial y'} + G_z^{qe}(r, r') \frac{\partial J_z}{\partial z'} \right] dS' \quad (4)$$

$$F(r) = \epsilon_0 \int_S \bar{G}_F(r, r') \cdot M(r') dS' \quad (5)$$

$$V_m(r) = \frac{-1}{j\omega\mu_0} \int_S \left[G_x^{qm}(r, r') \frac{\partial J_x}{\partial x'} + G_y^{qm}(r, r') \frac{\partial J_y}{\partial y'} + G_z^{qm}(r, r') \frac{\partial J_z}{\partial z'} \right] dS' \quad (6)$$

where r and r' represent the position vectors for the observation point and the increment area dS' , respectively, and μ_0 and ϵ_0 denote the free-space permeability and permittivity, respectively.

2.1. The Half-Space Green's Function

The half-space Green's function can be impressed by the vector and scalar potentials, and the vector potential is not uniquely specified. In this paper, we express it as follows [14]:

$$\bar{\mathbf{G}}_A = (\hat{\mathbf{x}}\hat{\mathbf{x}} + \hat{\mathbf{y}}\hat{\mathbf{y}}) G_A^{xx} + \hat{\mathbf{z}}\hat{\mathbf{x}}G_A^{zx} + \hat{\mathbf{z}}\hat{\mathbf{y}}G_A^{zy} + \hat{\mathbf{z}}\hat{\mathbf{z}}G_A^{zz} \quad (7)$$

$$\bar{\mathbf{G}}_F = (\hat{\mathbf{x}}\hat{\mathbf{x}} + \hat{\mathbf{y}}\hat{\mathbf{y}}) G_F^{xx} + \hat{\mathbf{z}}\hat{\mathbf{x}}G_F^{zx} + \hat{\mathbf{z}}\hat{\mathbf{y}}G_F^{zy} + \hat{\mathbf{z}}\hat{\mathbf{z}}G_F^{zz} \quad (8)$$

with

$$G_A^{yy} = G_A^{xx}, \quad G_y^{qe} = G_x^{qe} \quad (9)$$

$$G_F^{yy} = G_F^{xx}, \quad G_F^{qe} = G_F^{qe} \quad (10)$$

where G_A^{xx} , G_A^{zx} , G_A^{zy} and G_A^{zz} denote the spatial domain half-space Green's function for the electric vector potentials; G_x^{qe} , G_y^{qe} and G_z^{qe} denote the spatial domain half-space Green's function for the electric scalar potentials; G_F^{xx} , G_F^{zx} , G_F^{zy} and G_F^{zz} denote the spatial domain half-space Green's function for the magnetic vector potentials; G_x^{qm} , G_y^{qm} , G_z^{qm} denote the spatial domain half-space Green's function for the magnetic scalar potentials. The detailed accounts of derivations are beyond the scope of this paper. Please see [15, 16] for a full derivation. The main formulations are presented as follows:

$$G_A^{xx} = \frac{\mu}{4\pi} \left(-\frac{j}{2}\right) \int_{-\infty}^{+\infty} \frac{k_\rho}{k_z} H_0^{(2)}(k_\rho \rho) \left[e^{-jk_z|z-z'|} + A_h^e e^{-jk_z z} \right] dk_\rho \quad (11)$$

$$G_A^{zx} = -\frac{\mu}{8\pi} \frac{\partial}{\partial x} \int_{-\infty}^{+\infty} \frac{k_\rho}{k_z} H_0^{(2)}(k_\rho \rho) \left[\frac{k_z}{k_\rho^2} (C_h^e - A_h^e) e^{-jk_z z} \right] dk_\rho \quad (12)$$

$$G_x^{qe} = -\frac{j}{8\pi\epsilon} \int_{-\infty}^{+\infty} \frac{k_\rho}{k_z} H_0^{(2)}(k_\rho \rho) \left[e^{-jk_z|z-z'|} + \frac{k^2 A_h^e + k_z^2 C_h^e}{k_\rho^2} e^{jk_z z} \right] dk_\rho \quad (13)$$

$$G_A^{zz} = \frac{\mu}{4\pi} \left(-\frac{j}{2}\right) \int_{-\infty}^{+\infty} \frac{k_\rho}{k_z} H_0^{(2)}(k_\rho \rho) \left[e^{-jk_z|z-z'|} + A_v^e e^{-jk_z z} \right] dk_\rho \quad (14)$$

$$G_z^{qe} = \frac{1}{4\pi\epsilon} \left(-\frac{j}{2}\right) \int_{-\infty}^{+\infty} \frac{k_\rho}{k_z} H_0^{(2)}(k_\rho \rho) \left[e^{-jk_z|z-z'|} + C_v^e e^{-jk_z z} \right] dk_\rho \quad (15)$$

with

$$A_h^e = R_{\text{TE}} e^{-jk_z z'}, \quad C_h^e = -R_{\text{TM}} e^{-jk_z z'} \quad (16)$$

$$A_v^e = R_{\text{TE}} e^{-jk_z z'}, \quad C_v^e = -R_{\text{TM}} e^{-jk_z z'} \quad (17)$$

where R_{TE} , R_{TM} are Fresnel reflection coefficient, k_x, k_y, k_z are the wavenumbers of x, y, z in the global coordinate system, respectively. $k_\rho^2 = k_x^2 + k_y^2$, $H_0^{(2)}(\cdot)$ is the Hankel function of the second kind.

In order to obtain the half-space Green's function, we compute the vector and the scalar potentials by the discrete complex image method. Through the Sommerfeld identity, a closed-form spatial Green's function of a few terms is found from the quasi-dynamic images, the complex images, and the surface waves [15, 16].

To obtain the spatial-domain Green's functions in closed forms, it would be useful to give the definition of the spatial-domain Green's function

$$G = \frac{1}{4\pi} \int_{SIP} \tilde{G}(k_\rho) H_0^{(2)}(k_\rho) k_\rho dk_\rho \quad (18)$$

where, G and \tilde{G} are the Green's function in the spatial and spectral domains, respectively. $H_0^{(2)}$ is the hankel function of the second kind and SIP is the Sommerfeld integration path.

To alleviate the necessity of investigating the spectral-domain Green's functions in advance and the difficulties caused by the trade-off between the sampling range and the sampling period, the approximation is performed in two levels. The first part of approximation is performed along the path $Cap1$ while the second part is done along the path $Cap2$, as show in Fig. 2. The main formulations may refer to [17, 18].

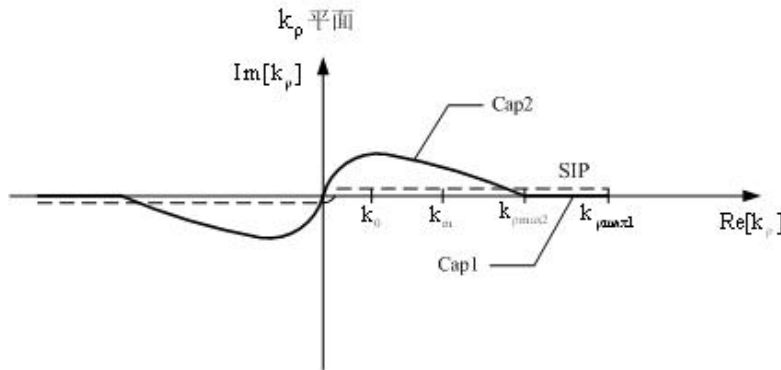


Figure 2. The path **Cap1** and **Cap2** used in two-level approximation.

2.2. Graphical Processing in Half-Space

As stated before, the half-space Green's function has been introduced into the high-frequency method to compute the electrically large targets in the half space. Combined with the graphical electromagnetic computing method, the geometry information of each illuminated

plane is obtained by reading the colors and depths of each pixel, and the shadow regions are eliminated by displaying lists technology of OpenGL to rebuild the target.

GRECO method was introduced by Rius in 1993 [3]. It can be integrated with CAD geometric modeling package and high-frequency theory for RCS prediction. By the hardware graphics accelerator, hidden surfaces of the image have been previously removed.

To the fixed screen, the calculation complexity is not varied with the complexity and dimension of targets. The targets are rebuilt by displaying lists technology of OpenGL and hidden surfaces of the image have been previously removed by the hardware graphics accelerator. Then we make use of the resolution to disperse the curve face into pixels that satisfy the requirement of the electromagnetic calculation, meanwhile, the scene is rendered using the Phong local illumination model [3]. For three light sources of purely green, red and blue colors, respectively, located over each one of the three coordinate axis, the three color components (Red, Blue, Green) for this pixel are equal to the (n_x, n_y, n_z) components of the unit normal to surface. Meanwhile, the depths of each pixel are obtained in the same way. The depth of each pixel is equal to the distance between the observer and each surface element.

Obtaining the unit normal and depth of each illuminated pixel of the target, the radar cross section of homogeneous dielectric object can be exactly calculated with the half-space physical optics method for homogeneous dielectric object below.

2.3. The Half-Space Physical Optics Method for Dielectric Object

The half-space Green's function has been introduced into the physical optics method to account for the scattering of homogeneous dielectric object with electrically large size in half space. Combined with the GRECO method, the shadow regions have been previously removed by the hardware graphics accelerator, and the geometry information is obtained by reading the color and depths of each pixel.

The total field backscattered field of the targets is summed up from the contributions of all visible facets while for the hidden parts the total fields are assumed to be zero. Based on the physical optics method for dielectric object, the surface currents (J_s, M_s) may be expressed as [19]:

$$\begin{aligned}
 J_s = & \frac{1}{z} \left(-e_{\perp}^i (i \cdot n) - R_{\perp} e_{\perp}^s (s \cdot n) \right) E_{\perp} \\
 & + \frac{1}{z} \left((n \times e_{\perp}^i) + R_{//} (n \times e_{\perp}^s) \right) E_{//}
 \end{aligned} \quad (19)$$

$$M_s = \left((e_{\perp}^i \times n) + R_{\perp} (e_{\perp}^s \times n) \right) E_{\perp} - \left(e_{\perp}^i (n \cdot i) + R_{//} (n \cdot s) \right) E_{//} \quad (20)$$

with [20]

$$R_{\perp} = \frac{\eta_2 \cos \theta_i - \eta_1 \cos \theta_t}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t} \quad (21)$$

$$R_{//} = \frac{\eta_1 \cos \theta_i - \eta_2 \cos \theta_t}{\eta_1 \cos \theta_i + \eta_2 \cos \theta_t} \quad (22)$$

in which $R_{\perp}, R_{//}$ are the Fresnel reflection coefficients, $\eta_i = \sqrt{u_i/\varepsilon_i}$ is the wave impedance for medium i ($i = 1, 3$), respectively.

To apply the far-field approximation, that is, the observation point is located far enough from the scattering object, we could obtain

$$\nabla \Phi \approx jk \hat{s} \Phi \quad (23)$$

and substitute (3), (4), (5), (23) into the scattered field (1), which leads to the equation:

$$E^s(\vec{r}) = -j\omega\mu_0 \int_S \bar{G}_A(r, r') \cdot J(r') dS' + \frac{k \cdot \hat{s}}{\omega \cdot \varepsilon_0} \int_S \left[G_x^{qe} \frac{\partial}{\partial x} J(\vec{r}') + G_y^{qe} \frac{\partial}{\partial y} J(\vec{r}') + G_z^{qe} \frac{\partial}{\partial z} J(\vec{r}') \right] d\vec{r}' - j\omega\varepsilon_0 \int_S \bar{G}_F(r, r') \cdot M(r') dS' \quad (24)$$

Equation (24) is then calculated for every illuminated facet and the complex RCS due to scattering from every illuminated facet is calculated as [21]:

$$\sqrt{\sigma} = \lim_{R \rightarrow \infty} 2\sqrt{\pi}R \frac{E_s \cdot \hat{e}_r}{E_o} \exp(jkR) \quad (25)$$

where $\exp(jkR)$ is a phase term that has been introduced into the equation in order to account for the facet location with respect to the global coordinate system, Here R is the position vector for the facet's reference vertex with respect to the global coordinate system.

3. RESULTS

When the object is calculated with the half-space physical optics method for dielectric objects from the simple targets to the complex

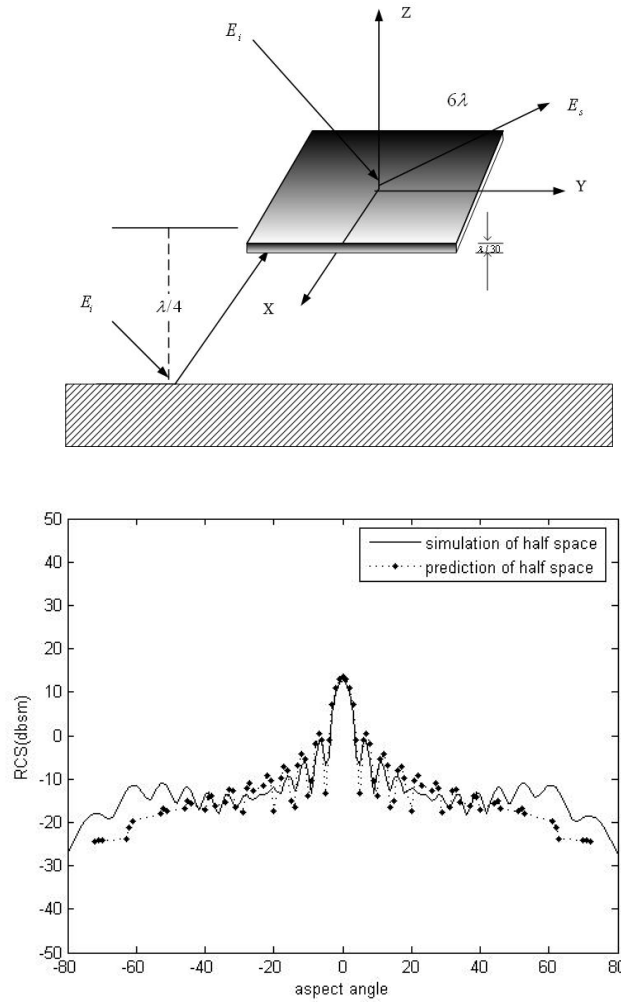


Figure 3. $6\text{wave} \times 6\text{wave} \times 1/30\text{wave}$ plane comparison of prediction and simulation of software in half space, frequency 10 GHz.

ones in half space, CPU time for RCS prediction is spent only on the electromagnetic part of the computation while the more time-consuming geometric model manipulations are left to the graphics hardware. Regardless the targets are the simple or complex ones, the scattering field of illuminated each pixel is calculated in the same way. Combined with the half-space Green's function, the high-frequency approximations to RCS prediction is easily computed from

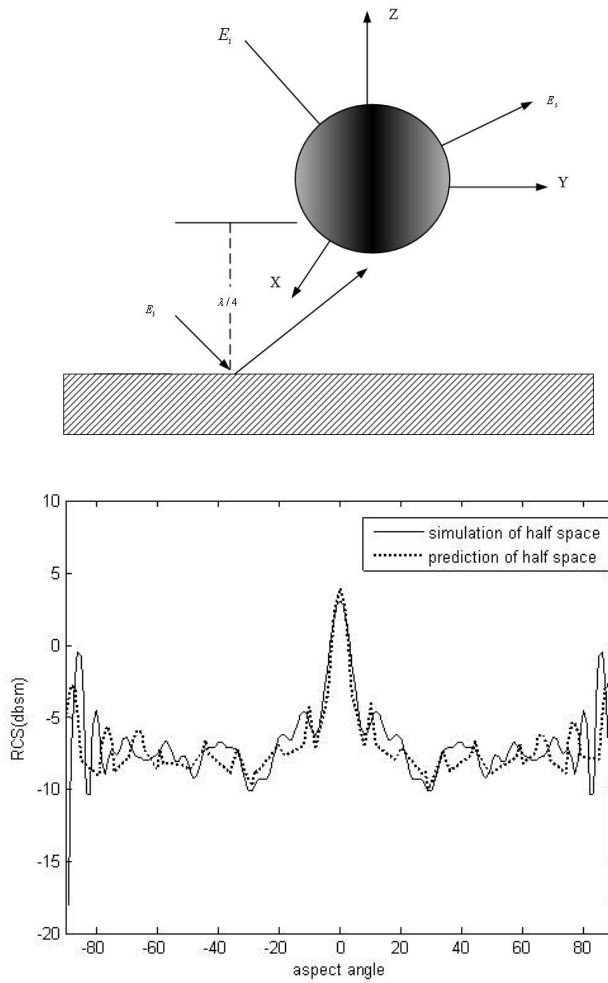


Figure 4. Dielectric sphere comparison of prediction and simulation of software in half space, frequency 10 GHz.

the knowledge of the unit normal at the illuminated surfaces of the targets in half space.

If consider the complex electrically large size ($L \gg \lambda$) targets in half space, we first compare with the results of the calculation and the simulation of the simple electrically large size targets in half space, such as plane, sphere, and apply this algorithm to the complex ones in half space. To ensure the accuracy of the derivated equation of the half-space physical optics method for dielectric object, we give two

examples below.

Figure 3 shows a dielectric flat plane ($6\lambda \times 6\lambda \times \frac{\lambda}{30}$) is placed $\lambda/4$ above the soil. Here, λ represents the wavelength in the vacuum. The dielectric plane has a relative dielectric permittivity $\epsilon_r = 4.0$, relative magnetic permeability $\mu_r = 1.0$. The soil has a relative dielectric permittivity $\epsilon_r = 4.0$, relative magnetic permeability $\mu_r = 1.0$.

There is good fit between the prediction of the derivated equation and the simulation of ANSOFT HFSS software, except for the flash produced by the leading edge of plate in total prediction. The reason for this discrepancy is caused by the currents that are induced on facet's free edges.

Figure 4 shows a dielectric sphere is placed $\lambda/4$ above the soil. Here, the sphere radius is 5λ and the soil has a relative dielectric permittivity $\epsilon_r = 4.0$, relative magnetic permeability $\mu_r = 1.0$. The dielectric sphere has a relative dielectric permittivity $\epsilon_r = 4.0$, relative magnetic permeability $\mu_r = 1.0$.

It indicates the predictions are in good agreement with the simulation of software, which proves that the half-space Green's function applied the physical optics method for dielectric object is accuracy for the targets with large electric dimensions in half space.

4. CONCLUSION

In this paper, we have presented an accurate and powerful approach for calculating the homogeneous dielectric object with electrically large size in half space. In order to consider the scattering fields of the targets in half space, the half-space physical optics method has been deduced by introducing the half-space Green's function into the conventional physical optics method (PO). Combined with the graphical-electromagnetic computing method to read the geometry information of all visible facets, the equivalent currents and the reflection coefficients were utilized to account of the homogenous dielectric objects with half-space physical optics method in half space. The numerical results have shown that this method is efficient and accurate.

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